

⇒ an coherent electronically excited state is a dynamic state that oscillates between the ground & excited state.

↳ Pump → excite the sys. into a coherent superposition. pump - probe
 ↳ probe measure how sys evolves ⇒ the detected signal oscillates in time.

1-5. spin precession. $(|\alpha, t_0; t\rangle)$

\hat{H} for spin $1/2$ sys.
 ↓
 (Energy)

magnetic energy $E = -\vec{\mu} \cdot \vec{B}$
↓
 magnetic dipole moment.

$\vec{\mu} = \gamma \cdot \vec{S}$
↓
 $+\frac{\hbar}{2}$: Angular momentum of particle.
↓
 gyromagnetic ratio

classic: related to particle's charge & mass.

by $\gamma = \frac{q\hbar}{2mc}$ g-factor (-2)
 more quantum: for e^- $\gamma = -\frac{e\hbar}{2meC} \cdot g_e$ (Landé)

$$\Rightarrow E = -\vec{\mu} \cdot \vec{B} = -\gamma \cdot \vec{S} \cdot \vec{B}$$

$$\Rightarrow \hat{H} = -\gamma \cdot \hat{S} \cdot \vec{B} \rightarrow (B_x, B_y, B_z)$$

$$\downarrow$$

$$(\hat{S}_x, \hat{S}_y, \hat{S}_z)$$

since $\hat{S}_x = \frac{\hbar}{2} \hat{\sigma}_x$, $\hat{S}_y = \frac{\hbar}{2} \hat{\sigma}_y$, $\hat{S}_z = \frac{\hbar}{2} \hat{\sigma}_z$

$$\Rightarrow \hat{H} = -\frac{\hbar}{2} \gamma \cdot \hat{\sigma} \cdot \vec{B}$$

↓
($\hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z$)

$$\Rightarrow \hat{H} = -\frac{\hbar}{2} \cdot \gamma \cdot (\hat{\sigma}_x B_x + \hat{\sigma}_y B_y + \hat{\sigma}_z B_z)$$

$$= -\frac{\hbar}{2} \cdot \gamma \cdot \left[\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} B_x + \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} B_y + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} B_z \right]$$

$$= -\frac{\hbar}{2} \cdot \gamma \cdot \begin{pmatrix} B_z & B_x - iB_y \\ B_x + iB_y & -B_z \end{pmatrix}$$

Now, if \vec{B} is static ($\frac{\partial \vec{B}}{\partial t} = 0$)

\vec{B} is only along $z \Rightarrow B_z$.

$$\Rightarrow \hat{H} = -\frac{\hbar}{2} \cdot \gamma \hat{\sigma}_z \cdot B_z$$

$$\Rightarrow [\hat{H}, \hat{\sigma}_z] = 0$$

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix} \& \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

↑

also, $\hat{\sigma}_z$ (\hat{S}_z) has eigen base kets. $|z; \pm\rangle$

Here, $\hat{\sigma}_z$ or \hat{S}_z is \hat{A} . $|z; \pm\rangle$ is $\{|a'\rangle\}$

suppose $|\alpha; t_0\rangle = C_+ |z; +\rangle + C_- |z; -\rangle$

what is $|\alpha, t_0; t\rangle = ?$ ($= \hat{U} |\alpha; t_0\rangle$)

$$\Rightarrow \hat{U}(t, t_0) = e^{-\frac{i}{\hbar} \hat{H} t}$$

$$= e^{-\frac{i}{\hbar} \left(-\frac{\hbar}{2} \cdot \gamma \hat{\sigma}_z \cdot B_z\right) t}$$

$$= e^{\frac{i}{\hbar} \gamma \hat{\sigma}_z B_z t} = e^{\frac{i}{\hbar} \gamma \hat{S}_z \cdot B_z \cdot t}$$

$$\hat{H} \cdot |\alpha; +\rangle = E_{\alpha+} |\alpha; +\rangle \quad \hat{H} \cdot |\alpha; -\rangle = E_{\alpha-} |\alpha; -\rangle$$

$$-\gamma \cdot \hat{S}_z \cdot B_z |\alpha; +\rangle = -\gamma \cdot B_z \cdot \frac{\hbar}{2} |\alpha; +\rangle$$

\downarrow
 $E_{\alpha+} = - \left(\frac{-e}{2m_0 c} \cdot g_e \right) \cdot B_z \cdot \frac{\hbar}{2}$
 \uparrow
 $= \frac{B_z e}{2m_0 c} \cdot \hbar$

likewise $E_{\alpha-} = - \frac{B_z e}{2m_0 c} \cdot \hbar$

$$\Rightarrow |\alpha; t_0; t\rangle = \sum_{\alpha'} C_{\alpha'} e^{-\frac{i}{\hbar} E_{\alpha'} t} |\alpha'\rangle$$

$$\Rightarrow \left(e^{-\frac{i}{\hbar} \frac{B_z e}{2m_0 c} \cdot \hbar \cdot t} \cdot C_+ |\alpha; +\rangle \right) + \left(e^{-\frac{i}{\hbar} \frac{-B_z e}{2m_0 c} \cdot \hbar \cdot t} \cdot C_- |\alpha; -\rangle \right)$$

$$= C_+ e^{-\frac{i\omega t}{2}} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + C_- e^{\frac{i\omega t}{2}} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\left(\omega = \frac{B_z e}{m_0 c} \right)$$

$$= \begin{pmatrix} C_+ e^{-i\omega t/2} \\ C_- e^{i\omega t/2} \end{pmatrix}$$

if $t=0 \Rightarrow |\alpha; t_0\rangle = \begin{pmatrix} C_+ \\ C_- \end{pmatrix}$

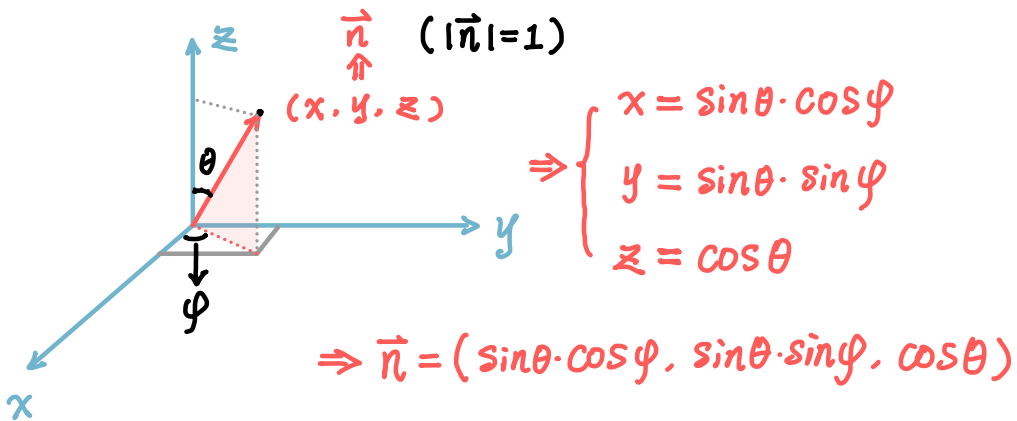
if $C_- = 0 \Rightarrow |\alpha; t_0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |\alpha; +\rangle$

⊗

NOTE: This is a very general case.
 Because $|x:\pm\rangle$ & $|y:\pm\rangle$ and other arbitrary forms can all be constructed by $|z:\pm\rangle$ based on C_+ & C_- .

Now, assuming we have an initial spin.

$$C_+ |z: +\rangle + C_- |z: -\rangle \quad \text{or} \quad \begin{pmatrix} C_+ \\ C_- \end{pmatrix}$$



This initial state $\begin{pmatrix} C_+ \\ C_- \end{pmatrix}$

must be the eigen state of $\hat{\sigma} \cdot \vec{n}$

w/ eigenvalue of unit. $= (\hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z)(x, y, z)$

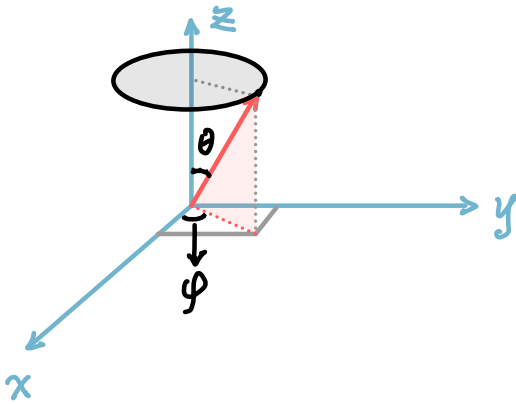
$$\Rightarrow \hat{\sigma} \cdot \vec{n} \begin{pmatrix} C_+ \\ C_- \end{pmatrix} = 1 \cdot \begin{pmatrix} C_+ \\ C_- \end{pmatrix}$$

$$\Rightarrow \text{solve } \begin{pmatrix} C_+ \\ C_- \end{pmatrix} = \begin{pmatrix} e^{-i\varphi/2} \cos(\theta/2) \\ e^{i\varphi/2} \sin(\theta/2) \end{pmatrix} \begin{matrix} \rightarrow C_+ \\ \rightarrow C_- \end{matrix}$$

NOW $\Rightarrow |\alpha, t_0; t\rangle$

$$= \begin{pmatrix} C_+ e^{-i\omega t/2} \\ C_- e^{i\omega t/2} \end{pmatrix}$$

$$= \begin{pmatrix} e^{\frac{-i(\omega t + \varphi)}{2}} \cos(\theta/2) \\ e^{\frac{i(\omega t + \varphi)}{2}} \sin(\theta/2) \end{pmatrix}$$



along z $\langle \hat{S}_z \rangle$

$$= \frac{\hbar}{2} |\langle z; + | \alpha, t_0; t \rangle|^2 + (-\frac{\hbar}{2}) |\langle z; - | \alpha, t_0; t \rangle|^2$$

$$= \frac{\hbar}{2} \cdot \cos\theta$$

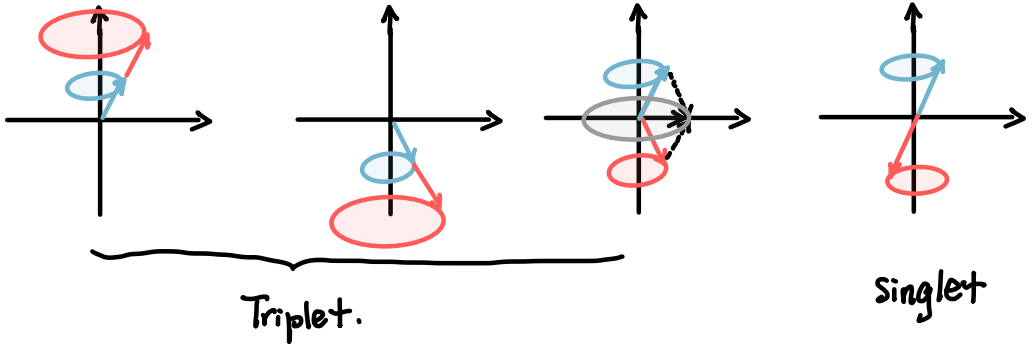
*. However along x & y

$\langle \hat{S}_x \rangle$ & $\langle \hat{S}_y \rangle$ will be related to \cos or \sin
 $(\omega t + \varphi)$

Triplet & singlet

($\uparrow\uparrow$)

($\uparrow\downarrow$)



1.6 Correlation Amplitude.

$$\begin{aligned} \text{Definition: } c(t) &\equiv \langle \underline{\alpha} | \alpha, t; t \rangle \\ &= \langle \alpha | \hat{u}(t, 0) | \alpha \rangle \end{aligned}$$

How much $|\alpha, t; t\rangle$ resemble the initial state $|\alpha\rangle$

$|c(t)|$ quantifies the similarity.

\Rightarrow if $|\alpha\rangle = |\alpha'\rangle$

$$\begin{aligned} \Rightarrow c(t) &= \langle \alpha' | \alpha', t; t \rangle = \langle \alpha' | e^{-\frac{i}{\hbar} E_{\alpha'} t} | \alpha' \rangle \\ &= e^{-\frac{i}{\hbar} E_{\alpha'} t} \end{aligned}$$

$\Rightarrow |c(t)| = 1 \Rightarrow$ stationary state.